# Potential Surfaces for Unimolecular and Bimolecular Gas Phase Reactions of $\mathbf{B H}_{\boldsymbol{m}} \mathbf{C l}_{\boldsymbol{n}}$ Calculated at the G2 Level of Theory 

H. Bernhard Schlegel* and Anwar G. Baboul<br>Department of Chemistry, Wayne State University, Detroit, Michigan 48202<br>Stephen J. Harris<br>Physical Chemistry, General Motors Research \& Development, Warren, Michigan 48090-9055

Received: February 5, 1996; In Final Form: March 5, $1996^{\otimes}$


#### Abstract

Transition structures and reaction paths for the $\mathrm{BH}_{m} \mathrm{Cl}_{n}$ system have been characterized at the MP2(full)/6$31 \mathrm{G}(\mathrm{d})$ level of theory; heats of reaction and barriers have been computed at the G2 level of theory. Calculations show that the insertion reactions of $\mathrm{BH}_{m} \mathrm{Cl}_{n}$ into $\mathrm{H}_{2}$ and $\mathrm{HCl}(m+n=0,1,2)$ occur by highly distorted, non-least motion transition states, with barriers that increase with chlorine substitution on boron. Hydrogen abstraction from BH and $\mathrm{BH}_{3-n} \mathrm{Cl}_{n}$ by chlorine proceeds with little or no barrier by a linear transition state. By contrast, $\mathrm{BH}_{3-n} \mathrm{Cl}_{n}+\mathrm{H} \rightarrow \mathrm{BH}_{2-n} \mathrm{Cl}_{n}+\mathrm{H}_{2}$ occurs by an addition-elimination path via a tetracoordinate intermediate. Likewise, $\mathrm{BH}_{2-n} \mathrm{Cl}_{n}+\mathrm{X} \rightarrow \mathrm{BH}_{1-n} \mathrm{Cl}_{n}+\mathrm{HX}$ or $\mathrm{BH}_{2-n} \mathrm{Cl}_{n-1}+\mathrm{XCl}(\mathrm{X}=\mathrm{H}$, $\mathrm{Cl})$ are also addition-elimination processes. Abstraction of Cl from BCl and $\mathrm{BH}_{3-n} \mathrm{Cl}_{n}$ by H is endothermic and occurs by a bent transition state.


## Introduction

Boron compounds such as cubic boron nitride, boron carbide, and titanium diboride are among the hardest and most abrasionresistant materials. Like diamond, these materials also have low densities, high melting points, and considerable mechanical strength and are chemically inert. As a result, they can make excellent wear- and corrosion-resistant coatings. ${ }^{1}$ Boroncontaining films can be grown using chemical vapor deposition (CVD). Because of the extreme toxicity associated with boron hydrides, $\mathrm{BCl}_{3}$ has been a commonly used boron precursor in CVD. Although there have been a number of publications examining the chemistry and kinetics of CVD film formation from $\mathrm{BCl}_{3}$, the validity of these analyses is uncertain because of a lack of accurate, relevant experimental or theoretical data on the decomposition of $\mathrm{BCl}_{3}$. For example, one of the most careful and detailed studies of boron carbide formation with CVD techniques modeled the chemistry and the fluid mechanics in a $\mathrm{BCl}_{3} / \mathrm{CH}_{4} / \mathrm{H}_{2}$ stagnation flow system with just two gas phase and two gas-surface reactions ${ }^{2}$

$$
\begin{gather*}
\mathrm{BCl}_{3}+\mathrm{H}_{2} \rightarrow \mathrm{BHCl}_{2}+\mathrm{HCl} \\
\mathrm{BHCl}_{2} \rightarrow \mathrm{BCl}+\mathrm{HCl} \\
\mathrm{BCl}+{ }^{1} /{ }_{2} \mathrm{H}_{2} \rightarrow \mathrm{~B}_{\mathrm{s}}+\mathrm{HCl} \\
\mathrm{CH}_{4} \rightarrow \mathrm{C}_{\mathrm{s}}+2 \mathrm{H}_{2} \tag{1}
\end{gather*}
$$

However, because of a lack of experimental or theoretical data for these species, there is in fact no evidence that any of these global reactions play a role in the formation of boron carbide. CVD reaction schemes leading to formation of boron nitride from $\mathrm{BCl}_{3}$ are of necessity similarly crude.

We have recently used quantum calculations at the G2 level of theory to provide a self-consistent set of heats of formation

[^0]

Figure 1. Potential reactions interconverting species in the $\mathrm{BH}_{m} \mathrm{Cl}_{n}$ system.
and vibrational frequencies for $\mathrm{BH}_{m} \mathrm{Cl}_{n}(0 \leq m, n \leq 3) .{ }^{3}$ These data can form part of the foundation for accurate modeling of $\mathrm{BCl}_{3}$ decomposition chemistry in the presence of $\mathrm{H}_{2}$, but kinetic information is also required. Figure 1 summarizes the unimolecular and bimolecular reactions that could interconvert the different species in the $\mathrm{BH}_{m} \mathrm{Cl}_{n}$ system. Experimental rate data are available for $\mathrm{BH}+\mathrm{H}_{2}^{4,5}$ and $\mathrm{BCl}_{3}+\mathrm{H}^{6}$ but not for most of the other reactions in Figure 1. The $\mathrm{BH}+\mathrm{H}_{2} \rightarrow \mathrm{BH}_{3}$ reaction has also been studied by high-level electronic structure methods. ${ }^{5}$ The reaction path and potential energy surface were obtained by CASSCF and MRCI calculations with large, correlation consistent basis sets. The theoretical rate constant and activation energy are in excellent agreement with experiment. The rates for $\mathrm{BCl}+\mathrm{H}, \mathrm{BCl}+\mathrm{Cl}, \mathrm{BH}+\mathrm{H}, \mathrm{BH}+\mathrm{Cl}$, and $\mathrm{B}+\mathrm{Cl}_{2}$ have also been examined but only using empirical BEBO surfaces. ${ }^{7,8}$ There are no systematic, high-level studies that treat the reactions of $\mathrm{BH}_{m} \mathrm{Cl}_{n}$ in a uniform and consistent manner. Theoretical methods such as the G2 level of theory are capable of calculating energetics of small gas phase species to within $1-2 \mathrm{kcal} / \mathrm{mol} .^{9-11}$ Previously, we have used this level of theory to calculate the heats of formation of $\mathrm{BH}_{m} \mathrm{Cl}_{n}{ }^{3}$ In this paper, we use the same approach to investigate the transition structures and reaction barriers in the $\mathrm{BH}_{m} \mathrm{Cl}_{n}$ system.

TABLE 1: Total Energies (hartrees) for Reactants, Intermediates, and Transition States in the $\mathbf{B H}_{m} \mathbf{C l}_{n} \mathbf{S y s t e m}^{a}$

| species | sym | MP2(full)/6-31G(d) ${ }^{\text {b }}$ | $\mathrm{G} 2{ }^{\text {c }}$ |
| :---: | :---: | :---: | :---: |
| reactants and intermediates |  |  |  |
| BH | $C_{\infty v}$ | -25.175 870 | -25.233 699 |
| BCl | $C_{\infty}$ | -484.311686 | -484.474 419 |
| $\mathrm{BH}_{2}$ | $C_{2 v}$ | -25.808 445 | -25.856 344 |
| BHCl | $C_{5}$ | -484.911 069 | -485.067 923 |
| $\mathrm{BCl}_{2}$ | $C_{2 v}$ | -944.008 108 | -944.274 167 |
| $\mathrm{BH}_{3}$ | $D_{3 h}$ | -26.468 568 | -26.523 328 |
| $\mathrm{BH}_{2} \mathrm{Cl}$ | $C_{2 v}$ | -485.570 245 | -485.733 340 |
| $\mathrm{BHCl}_{2}$ | $C_{2 v}$ | -944.668 094 | -944.940 296 |
| $\mathrm{BCl}_{3}$ | $D_{3 h}$ | -1403.759581 | -1404.141597 |
| $\mathrm{BH}_{2}-\mathrm{H}_{2}$ | $C_{2 v}$ | -26.977 691 | -27.046 430 |
| $\mathrm{BHCl}-\mathrm{H}_{2}$ | $C_{s}$ | -486.069 302 | -486.244 956 |
| $\mathrm{BCl}_{2}-\mathrm{H}_{2}$ | $C_{2 v}$ | -945.162 373 | -945.445738 |
| insertions transition states |  |  |  |
| $\mathrm{B}+\mathrm{H}_{2} \rightarrow \mathrm{BH}_{2}$ | $C_{s}$ | -25.658 867 | -25.737081 |
| $\mathrm{BH}+\mathrm{H}_{2} \rightarrow \mathrm{BH}_{3}$ | $C_{s}$ | -26.309 448 | -26.392 132 |
| $\mathrm{BCl}+\mathrm{H}_{2} \rightarrow \mathrm{BH}_{2} \mathrm{Cl}$ | $C_{s}$ | -485.396 447 | -485.589 758 |
| $\mathrm{BH}_{2}+\mathrm{H}_{2} \rightarrow \mathrm{BH}_{2}-\mathrm{H}_{2}$ | $C_{s}$ | -26.951 144 | -27.021 788 |
| $\mathrm{BHCl}+\mathrm{H}_{2} \rightarrow \mathrm{BHCl}-\mathrm{H}_{2}$ | $C_{1}$ | -486.041 908 | -486.223 151 |
| $\mathrm{BCl}_{2}+\mathrm{H}_{2} \rightarrow \mathrm{BCl}_{2}-\mathrm{H}_{2}$ | $C_{s}$ | -945.128 278 | -945.421 062 |
| $\mathrm{B}+\mathrm{HCl} \rightarrow \mathrm{BHCl}$ | $C_{s}$ | -484.752 255 | -484.942865 |
| $\mathrm{BH}+\mathrm{HCl} \rightarrow \mathrm{BH}_{2} \mathrm{Cl}$ | $C_{s}$ | -485.384 834 | -485.576 476 |
| $\mathrm{BCl}+\mathrm{HCl} \rightarrow \mathrm{BHCl}_{2}$ | $C_{s}$ | -944.505 180 | -944.804 380 |
| hydrogen abstractions transition states |  |  |  |
| $\mathrm{BH}+\mathrm{H} \rightarrow \mathrm{B}+\mathrm{H}_{2}$ | $C_{\infty v}$ | -25.656 946 | -25.726 773 |
| $\mathrm{BH}_{3}+\mathrm{H} \rightarrow \mathrm{BH}_{2}+\mathrm{H}_{2}{ }^{\text {d }}$ | $C_{2 v}$ | -26.929 235 | -27.003 450 |
| $\mathrm{BH}_{2} \mathrm{Cl}+\mathrm{H} \rightarrow \mathrm{BHCl}+\mathrm{H}_{2}$ | $C_{s}$ | -486.032 256 | -486.214 703 |
| $\mathrm{BHCl}_{2}+\mathrm{H} \rightarrow \mathrm{BCl}_{2}+\mathrm{H}_{2}$ | $C_{2 v}$ | -945.130 842 | -945.421 631 |
| $\mathrm{BH}_{3}+\mathrm{Cl} \rightarrow \mathrm{BH}_{2}+\mathrm{HCl}$ | $C_{2 v}$ | -486.010 676 | -486.199 347 |
| $\mathrm{BH}_{2} \mathrm{Cl}+\mathrm{Cl} \rightarrow \mathrm{BHCl}+\mathrm{HCl}$ | $C_{s}$ | -945.133 354 | -945.411 155 |
| $\mathrm{BHCl}_{2}+\mathrm{Cl} \rightarrow \mathrm{BCl}_{2}+\mathrm{HCl}$ | $C_{2 v}$ | -1404.210 355 | -1404.616958 |
| chlorine abstractions transition states |  |  |  |
| $\mathrm{BCl}+\mathrm{H} \rightarrow \mathrm{B}+\mathrm{HCl}$ | $C_{s}$ | -484.740 574 | -484.933 030 |
| $\mathrm{BH}_{2} \mathrm{Cl}+\mathrm{H} \rightarrow \mathrm{BH}_{2}+\mathrm{HCl}$ | $C_{s}$ | -486.009 238 | -486.200 321 |
| $\mathrm{BHCl}_{2}+\mathrm{H} \rightarrow \mathrm{BHCl}+\mathrm{HCl}$ | $C_{1}$ | -945.100 220 | -945.400 346 |
| $\mathrm{BCl}_{3}+\mathrm{H} \rightarrow \mathrm{BCl}_{2}+\mathrm{HCl}$ | $C_{s}$ | -1404.188 613 | -1404.598971 |

${ }^{a}$ Additional MP2 ${ }^{b}$ and G2 ${ }^{c}$ energies: $\mathrm{H}-0.498$ 230, -0.500000 ; B $-24.562460,-24.60205 ; \mathrm{Cl}-459.562060,-459.676640 ; \mathrm{H}_{2}-1.144141$, -1.165863 ; $\mathrm{HCl}-460.202149,-460.339955 .^{b}$ Without ZPE. ${ }^{c}$ At 0 K , with ZPE using MP2(full)/6-31G(d) frequencies scaled by 0.9646 rather than HF/6-31G(d) frequences. ${ }^{d}$ Second-order saddle point

## Method

Molecular orbital calculations were carried out using the GAUSSIAN $94^{12}$ series of programs using a variety of basis sets of split valence quality or better with multiple polarization and diffuse functions ${ }^{13}$ and spin-unrestricted methods. Equilibrium geometries were optimized by Hartree-Fock and second-order Møller-Plesset perturbation theory (HF/6-31G(d) and MP2(full)/6-31G(d), respectively) using a quasi-Newton optimization method. ${ }^{14}$ Vibrational frequencies and zero-point energies were calculated at the HF/6-31G(d) and MP2(full)/6$31 \mathrm{G}(\mathrm{d})$ levels using analytical second derivatives. ${ }^{15,16}$ All transition states had only one imaginary frequency. Some representative transition states were characterized further by reaction path following ${ }^{17}$ to confirm the nature of the reactants and products of the reaction. Correlated energies were calculated by fourth-order Møller-Plesset perturbation theory ${ }^{18}$ (MP4SDTQ, frozen core) and by quadratic configuration interaction with perturbative correction for triple excitations ${ }^{19}$ ( $\mathrm{QCISD}(\mathrm{T})$, frozen core) with the MP2(full)/6-31G(d) optimized geometries. In the G2 method, ${ }^{9,10}$ the energy computed at MP4/ $6-311 \mathrm{G}(\mathrm{d}, \mathrm{p})$ is corrected for the effect of diffuse functions obtained at MP4/6-311+G(d,p) and MP2/6-311+G(3df,2p), for the effect of higher polarization functions obtained at MP4/6$311 \mathrm{G}(2 \mathrm{df}, \mathrm{p}$ ) and MP2/6-311+G(3df,2p), and for the effect of electron correlation beyond fourth-order obtained at QCISD-(T)/6-311G(d,p). Higher level corrections for deficiencies in the wave function are estimated empirically ${ }^{9,10}$ based on the number of paired and unpaired electrons.

$$
\begin{align*}
& E(\mathrm{G} 2)=E(\mathrm{MP} 4 / 6-311 \mathrm{G}(\mathrm{~d}, \mathrm{p}))+\Delta E(\text { diffuse })+ \\
& \quad \Delta E(\text { polarization })+\Delta E(\mathrm{QCI})+\Delta E(\mathrm{HLC})+\mathrm{ZPE} \tag{2}
\end{align*}
$$

The average absolute error of the additivity assumptions in the G2 level of theory is only $0.30 \mathrm{kcal} / \mathrm{mol} .{ }^{11}$ Because the geometry and frequencies of some transition states are sensitive to the level of theory, in the present work we use the zeropoint energy calculated at the MP2/6-31G(d) scaled by $0.9646^{20}$ for all structures (minima and transition states) rather than the HF/6-31G(d) zero-point energy scaled by $0.8929 .{ }^{21}$ At the G2 level of theory, the mean absolute error for 125 well-characterized atomization energies, ionization energies, electron affinities, and proton affinities is $1.3 \mathrm{kcal} / \mathrm{mol} .{ }^{10}$ The errors for barrier heights are probably similar or somewhat larger.

## Results and Discussion

The reactions in Figure 1 can be grouped into three basic categories: homolytic bond cleavage reactions, radical abstractions, and insertion (or 1,1-elimination) reactions. Some of the reactions in Figure 1 can also occur by an addition-elimination process. The reverse of the bond cleavage, radical addition reactions, should have little or no barrier. Hence, we did not search for a classical transition state on the potential energy surface for these reactions (however, a variational transition state would exist on the free energy surface).

The MP2 and G2 energies of various stationary points on the $\mathrm{BH}_{m} \mathrm{Cl}_{n}$ surface are listed in Table 1. Most of the openshell species have $S^{2}$ values of $0.75-0.85$, except the transition states for $\mathrm{B}+\mathrm{H}_{2}, \mathrm{~B}+\mathrm{HCl}$, and $\mathrm{BCl}+\mathrm{H}$, which are in the range $0.95-1.1$. The geometries optimized at the MP2(full)/ $6-31 G(d)$ level are given in Figures $2-8$. While this level of theory is usually quite satisfactory for equilibrium geometries, the geometries of some transition states may improve with larger basis sets and higher levels of theory. The harmonic vibrational

TABLE 2: Moments of Inertia and Frequencies for Reactions, Intermediates, and Transition States ${ }^{\boldsymbol{a}}$

| BH | 5.017 |
| :---: | :---: |
|  | 2448.4 |
| BCl | 87.854 |
|  | 880.2 |
| $\mathrm{BH}_{2}$ | 1.675, 8.183, 9.858 |
|  | 1080.2, 2688.3, 2847.9 |
| BHCl | 3.017, 101.764, 104.781 |
|  | 884.3, 969.9, 2744.9 |
| $\mathrm{BCl}_{2}$ | 21.187, 583.362, 604.549 |
|  | 299.7, 742.1, 1034.5 |
| $\mathrm{BH}_{3}$ | 7.667, 7.667, 15.333 |
|  | 1191.8, 1255.1, 1255.1, 2646.9, 2789.8, 2789.8 |
| $\mathrm{BH}_{2} \mathrm{Cl}$ | 7.872, 114.279, 122.151 |
|  | 868.5, 919.7, 1044.2, 1291.2, 2709.7, 2826.4 |
| $\mathrm{BHCl}_{2}$ | 37.659, 568.834, 606.493 |
|  | 308.6, 772.8, 824.3, 953.7, 1167.2, 2792.8 |
| $\mathrm{BCl}_{3}$ | 564.086, 564.086, 1128.172 |
|  | 272.6, 272.6, 479.9, 496.9, 1009.1, 1009.1 |
| $\mathrm{BH}_{2}-\mathrm{H}_{2}$ | 10.342, 13.238, 19.206 |
|  | 803.1, 858.1, 889.5, 1109.0, 1321.7, 2110.5, 2224.2, 2731.9, 2852.9 |
| $\mathrm{BHCl}-\mathrm{H}_{2}$ | 13.251, 130.938, 137.676 |
|  | 685.0, 777.2, 810.2, 855.3, 939.4, 1221.0, 1963.8, 2201.9, 2784.6 |
| $\mathrm{BCl}_{2}-\mathrm{H}_{2}$ | 50.878, 588.745, 631.568 |
|  | $289.5,491.8,740.5,821.0,829.5,955.4,1015.3,1914.8,2211.4$ |
| insertions transition states |  |
| $\mathrm{B}+\mathrm{H}_{2} \rightarrow \mathrm{BH}_{2}$ | 2.527, 13.286, 15.813 |
|  | 1100.4i, 698.3, 2195.7 |
| $\mathrm{BH}+\mathrm{H}_{2} \rightarrow \mathrm{BH}_{3}$ | 6.038, 15.536, 21.574 |
|  | 626.4i, 911.9, 1104.4, 1423.3, 2517.9, 3293.4 |
| $\mathrm{BCl}+\mathrm{H}_{2} \rightarrow \mathrm{BH}_{2} \mathrm{Cl}$ | 9.391, 114.631, 124.022 |
|  | 1446.5i, 669.1, 871.3, 935.0, 1450.0, 2509.1 |
| $\mathrm{BH}_{2}+\mathrm{H}_{2} \rightarrow \mathrm{BH}_{2}-\mathrm{H}_{2}$ | 10.748, 26.237, 31.970 |
|  | 388.5i, 310.7, 616.4, 688.2, 917.5, 1039.5, 2703.0, 2866.1, 4291.5 |
| $\mathrm{BHCl}+\mathrm{H}_{2} \rightarrow \mathrm{BHCl}-\mathrm{H}_{2}$ | 19.134, 132.513, 145.092 |
|  | 695.5i, 503.3, 589.5, 763.4, 845.2, 964.7, 1112.6, 2783.4, 3904.8 |
| $\mathrm{BCl}_{2}+\mathrm{H}_{2} \rightarrow \mathrm{BCl}_{2}-\mathrm{H}_{2}$ | 53.919, 599.387, 641.209 |
|  | 805.0i, 283.2, 459.1, 575.2, 702.7, 806.1, 1001.4, 1225.5, 3331.1 |
| $\mathrm{B}+\mathrm{HCl} \rightarrow \mathrm{BHCl}$ | 2.304, 184.119, 186.422 |
|  | 879.9i, 333.1, 1470.2 |
| $\mathrm{BH}+\mathrm{HCl} \rightarrow \mathrm{BH}_{2} \mathrm{Cl}$ | 4.972, 269.385, 274.357 |
|  | 487.5i, 225.6, 766.1, 780.4, 1432.5, 2624.7 |
| $\mathrm{BCl}+\mathrm{HCl} \rightarrow \mathrm{BHCl}_{2}$ | 41.635, 898.382, 940.018 |
|  | $737.2 i, 118.5,518.8,735.5,962.6,1572.1$ |
| hydrogen abstractions transition states |  |
| $\mathrm{BH}+\mathrm{H} \rightarrow \mathrm{B}+\mathrm{H}_{2}$ | 25.515, 25.515 |
|  | 2015.4i, 1020.5, 1020.5, 1160.6 |
| $\mathrm{BH}_{3}+\mathrm{H} \rightarrow \mathrm{BH}_{2}+\mathrm{H}_{2}{ }^{\text {b }}$ | 8.023, 30.092, 38.116 |
|  | 1749.3i, 386.7i, 389.7, 1122.0, 1180.3, 1207.5, 1602.9, 2688.4, 2832.1 |
| $\mathrm{BH}_{2} \mathrm{Cl}+\mathrm{H} \rightarrow \mathrm{BHCl}+\mathrm{H}_{2}$ | 23.267, 137.120, 160.387 |
|  | $1764.9 i, 88.2,236.3,883.5,956.8,1184.9,1223.4,1598.2,2750.4$ |
| $\mathrm{BHCl}_{2}+\mathrm{H} \rightarrow \mathrm{BCl}_{2}+\mathrm{H}_{2}$ | 73.465, 578.697, 652.163 |
|  | 1728.1i, 133.0, 220.3, 298.5, 707.6, 1008.8, 1158.9, 1219.4, 1629.9 |
| $\mathrm{BH}_{3}+\mathrm{Cl} \rightarrow \mathrm{BH}_{2}+\mathrm{HCl}$ | 8.398, 331.911, 340.308 |
|  | 568.8i, 229.3, 248.0, 817.6, 925.9, 958.8, 1093.4, 2718.6, 2895.5 |
| $\mathrm{BH}_{2} \mathrm{Cl}+\mathrm{Cl} \rightarrow \mathrm{BHCl}+\mathrm{HCl}$ | 57.702, 1073.396, 1131.098 |
|  | 514.5i, 83.5, 240.0, 769.5, 815.3, 900.6, 940.2, 1056.7, 2779.8 |
| $\mathrm{BHCl}_{2}+\mathrm{Cl} \rightarrow \mathrm{BCl}_{2}+\mathrm{HCl}$ | 588.839, 1190.769, 1779.608 |
|  | $435.9 i, 62.2,151.1,260.6,640.6,716.8,828.9,1043.9,1096.2$ |
| chlorine abstractions transition states |  |
| $\mathrm{BCl}+\mathrm{H} \rightarrow \mathrm{B}+\mathrm{HCl}$ | 4.354, 115.024, 119.378 |
|  | 1070.4i, 647.5, 1063.3 |
| $\mathrm{BH}_{2} \mathrm{Cl}+\mathrm{H} \rightarrow \mathrm{BH}_{2}+\mathrm{HCl}$ | 13.590, 136.187, 139.155 |
|  | 940.1i, 452.8, 623.0, 766.6, 869.1, 1070.1, 1337.9, 2713.3, 2861.8 |
| $\mathrm{BHCl}_{2}+\mathrm{H} \rightarrow \mathrm{BHCl}+\mathrm{HCl}$ | 51.607, 608.976, 649.492 |
|  | 846.8i, 238.7, 395.6, 542.5, 751.4, 871.2, 1037.9, 1455.3, 2790.0 |
| $\mathrm{BCl}_{3}+\mathrm{H} \rightarrow \mathrm{BCl}_{2}+\mathrm{HCl}$ | 591.793, 627.800, 1201.123 |
|  | 830.6i, 195.1, 209.6, 308.3, 388.2, 539.5, 766.8, 977.4, 1337.9 |

${ }^{a}$ For each entry: first line, moments of inertia (amu bohr ${ }^{2}$ ); second line, frequencies $\left(\mathrm{cm}^{-1}\right)$; based on ${ }^{1} \mathrm{H},{ }^{11} \mathrm{~B}$, and ${ }^{35} \mathrm{Cl}$; additional values: $\mathrm{H}_{2}$, $0.979,4532.4 ; \mathrm{HCl}, 5.733,3046.5 .{ }^{b}$ Second-order saddle point.
frequencies and moments of inertia calculated at the MP2(full)/ $6-31 \mathrm{G}(\mathrm{d})$ level are listed in Table 2. The calculated vibrational frequencies at this level of theory are typically ca. $6 \%$ too high ${ }^{20}$ when compared to observed frequencies because of basis set effects, the neglect of electron correlation, and vibrational
anharmonicity. Heats of reaction and barrier heights are given in Table 3. For some reactions, there may be weakly bound reactant-like clusters, allowing the transition states to be slightly lower than separated reactants. Note that recent experiments ${ }^{22,23}$ and calculations ${ }^{4,24}$ indicate that the heat of formation of gas

TABLE 3: Heats of Reaction and Barrier Heights in the $\mathbf{B H}_{m} \mathbf{C l}_{n}$ System $^{a}$

| species | heats of reaction |  | forward barrier |  | reverse barrier |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | MP2 | G2 | MP2 | G2 | MP2 | G2 |
| insertions transition states |  |  |  |  |  |  |
| $\mathrm{B}+\mathrm{H}_{2} \rightarrow \mathrm{BH}_{2}$ | -63.9 | -55.5 | 30.0 | 19.3 | 93.9 | 78.8 |
| $\mathrm{BH}+\mathrm{H}_{2} \rightarrow \mathrm{BH}_{3}$ | -93.2 | -77.7 | 6.6 | $4.6{ }^{\text {b }}$ | 99.8 | 82.3 |
| $\mathrm{BCl}+\mathrm{H}_{2} \rightarrow \mathrm{BH}_{2} \mathrm{Cl}$ | -71.8 | -58.4 | 37.3 | 31.7 | 109.1 | 90.1 |
| $\mathrm{BH}_{2}+\mathrm{H}_{2} \rightarrow \mathrm{BH}_{2}-\mathrm{H}_{2}$ | -15.8 | -15.2 | 0.9 | 0.4 | 16.6 | 15.6 |
| $\mathrm{BHCl}+\mathrm{H}_{2} \rightarrow \mathrm{BHCl}-\mathrm{H}_{2}$ | -8.8 | $-7.0$ | 8.4 | 6.7 | 17.2 | 13.7 |
| $\mathrm{BCl}_{2}+\mathrm{H}_{2} \rightarrow \mathrm{BCl}_{2}-\mathrm{H}_{2}$ | -6.4 | -3.6 | 15.0 | 12.0 | 21.4 | 15.6 |
| $\mathrm{B}+\mathrm{HCl} \rightarrow \mathrm{BHCl}$ | -91.9 | -79.0 | 7.8 | $-0.5$ | 99.7 | 78.5 |
| $\mathrm{BH}+\mathrm{HCl} \rightarrow \mathrm{BH}_{2} \mathrm{Cl}$ | -120.6 | $-100.2$ | -2.2 | -1.8 | 118.5 | 98.4 |
| $\mathrm{BCl}+\mathrm{HCl} \rightarrow \mathrm{BHCl}_{2}$ | 96.8 | -79.0 | 5.4 | 6.3 | 102.2 | 85.3 |
| hydrogen abstractions transition states |  |  |  |  |  |  |
| $\mathrm{BH}+\mathrm{H} \rightarrow \mathrm{B}+\mathrm{H}_{2}$ | -20.4 | -21.5 | 10.8 | 4.3 | 31.2 | 25.8 |
| $\mathrm{BH}_{3}+\mathrm{H} \rightarrow \mathrm{BH}_{2}+\mathrm{H}_{2}{ }^{\text {c }}$ | 8.9 | 0.7 | 23.6 | 12.5 | 14.7 | 11.8 |
| $\mathrm{BH}_{2} \mathrm{Cl}+\mathrm{H} \rightarrow \mathrm{BHCl}+\mathrm{H}_{2}$ | 8.3 | $-0.3$ | 22.7 | 11.7 | 14.4 | 12.0 |
| $\mathrm{BHCl}_{2}+\mathrm{H} \rightarrow \mathrm{BCl}_{2}+\mathrm{H}_{2}$ | 8.8 | 0.2 | 22.3 | 11.7 | 13.4 | 11.5 |
| $\mathrm{BH}+\mathrm{Cl} \rightarrow \mathrm{B}+\mathrm{HCl}$ | 12.6 | 2.3 | 12.5 | 0.4 | -0.1 | 1.9 |
| $\mathrm{BH}_{3}+\mathrm{Cl} \rightarrow \mathrm{BH}_{2}+\mathrm{HCl}$ | 12.6 | 2.3 | 12.5 | 0.4 | -0.1 | 1.9 |
| $\mathrm{BH}_{2} \mathrm{Cl}+\mathrm{Cl} \rightarrow \mathrm{BHCl}+\mathrm{HCl}$ | 12.0 | 1.3 | 11.9 | $-0.7$ | -0.1 | -2.0 |
| $\mathrm{BHCl}_{2}+\mathrm{Cl} \rightarrow \mathrm{BCl}_{2}+\mathrm{HCl}$ | 12.5 | 1.8 | 12.4 | 0.0 | -0.1 | $-1.8$ |
| chlorine abstractions transition states |  |  |  |  |  |  |
| $\mathrm{BCl}+\mathrm{H} \rightarrow \mathrm{B}+\mathrm{HCl}$ | 28.4 | 20.4 | 43.5 | 26.0 | 15.1 | 5.6 |
| $\mathrm{BH}_{2} \mathrm{Cl}+\mathrm{H} \rightarrow \mathrm{BH}_{2}+\mathrm{HCl}$ | 36.3 | 23.2 | 37.2 | 20.7 | 0.9 | $-2.5$ |
| $\mathrm{BHCl}_{2}+\mathrm{H} \rightarrow \mathrm{BHCl}+\mathrm{HCl}$ | 33.3 | 20.3 | 41.5 | 25.1 | 8.2 | 4.7 |
| $\mathrm{BCl}_{3}+\mathrm{H} \rightarrow \mathrm{BCl}_{2}+\mathrm{HCl}$ | 29.8 | 17.2 | 43.4 | 26.7 | 13.6 | 9.5 |

${ }^{a}$ At $0 \mathrm{~K}, \mathrm{MP} 2$ values without ZPE, G2 values include ZPE. ${ }^{b}$ Compare with ref 5, $2.9 \mathrm{kcal} / \mathrm{mol}$ calculated by MRCI/pVTZ+sp. ${ }^{c}$ Second-order saddle point.


Figure 2. Transition states for the insertion of $\mathrm{B}, \mathrm{BH}$, and BCl into $\mathrm{H}_{2}$ optimized at the MP2(full)/6-31G(d) level of theory.
phase boron atom reported in some standard reference works (e.g., $133.3 \mathrm{kcal} / \mathrm{mol}$ in ref 25 ) is ca. $3 \mathrm{kcal} / \mathrm{mol}$ too low. The experimental value of $136.2 \mathrm{kcal} / \mathrm{mol}$ determined by Storms and Mueller ${ }^{22}$ is supported by photoionization experiments ${ }^{23}$ and by high-level calculations of Ochterski, Petersson, and Wiberg ${ }^{24}$ and by our previous theoretical work. ${ }^{3}$

The transition states for insertion reactions of $\mathrm{B}, \mathrm{BH}$, and BCl into $\mathrm{H}_{2}$ are shown in Figure 2. All of these insertion reactions have a strongly distorted, non-least motion transition structure because of orbital symmetry considerations similar to carbene and silylene insertions. ${ }^{26}$ Like the carbene and silylene insertions, ${ }^{26}$ these reactions show a profound increase in the barrier height on halogen substitution $(4.6 \mathrm{kcal} / \mathrm{mol}$ for $\mathrm{BH}+$ $\mathrm{H}_{2}$ versus $31.7 \mathrm{kcal} / \mathrm{mol}$ for $\mathrm{BCl}+\mathrm{H}_{2}$ ). The $\mathrm{BH}+\mathrm{H}_{2} \rightarrow \mathrm{BH}_{3}$ reaction has been studied previously at higher levels of theory. The MP2/6-31G(d) optimized transition state is ca. $0.1 \AA$ earlier than the CASSCF/pVTZ-f structure; ${ }^{5}$ the MRCI/pVTZ+sp barrier also occurs significantly earlier along the reaction path. ${ }^{5}$ The barrier computed at the G2 level of theory is $1.7 \mathrm{kcal} / \mathrm{mol}$ higher than the $\mathrm{MRCI} / \mathrm{pVTZ}+\mathrm{sp}$ barrier. The latter is in excellent agreement with experiment. ${ }^{5}$

The transition states for the reaction of $\mathrm{B}, \mathrm{BH}$, and BCl with HCl shown in Figure 3 look more like abstractions, but the reaction path following confirms that these are insertion reactions. The reactants and transition states for $\mathrm{B}, \mathrm{BH}$, and


Figure 3. Transition states for the insertion of $\mathrm{B}, \mathrm{BH}$, and BCl into HCl optimized at the MP2(full)/6-31G(d) level of theory.

BCl plus HCl lie $20-40 \mathrm{kcal} / \mathrm{mol}$ below the abstraction products, $\mathrm{BH}, \mathrm{BH}_{2}$, and BHCl plus Cl , because of the strength of the $\mathrm{H}-\mathrm{Cl}$ bond being broken compared to the $\mathrm{B}-\mathrm{H}$ bond being formed. By contrast, the insertion reactions are exothermic by $100-120 \mathrm{kcal} / \mathrm{mol}$ because two bonds are formed but only one is broken. Somewhat surprisingly, the approach of the reagents appears to be governed by the HCl hydrogen interacting with the boron lone pair rather than the chlorine lone pair interacting with the empty boron p orbital.

Figure 4 shows the hydrogen abstraction transition structures for $\mathrm{BH}+\mathrm{H} \rightarrow \mathrm{B}+\mathrm{H}_{2}$ and $\mathrm{BH}_{3-\mathrm{n}} \mathrm{Cl}_{n}+\mathrm{H} \rightarrow \mathrm{BH}_{2-n} \mathrm{Cl}_{n}+\mathrm{H}_{2}$. No abstraction transition states could be found for $\mathrm{BH}_{2-n} \mathrm{Cl}_{n}+$ H , since the lowest energy processes are barrierless H additions to form $\mathrm{BH}_{3-n} \mathrm{Cl}_{n}$. Each of the four structures of the stationary points shown in Figure 4 has a linear $\mathrm{B}-\mathrm{H}-\mathrm{H}$ moiety, with a planar boron and partially broken/formed $\mathrm{B}-\mathrm{H}$ and $\mathrm{H}-\mathrm{H}$ bonds. Because of the weakness of the BH bond, $\mathrm{BH}+\mathrm{H} \rightarrow \mathrm{B}+\mathrm{H}_{2}$ is quite exothermic ( $21.5 \mathrm{kcal} / \mathrm{mol}$ ) and has a small barrier ( 4.3 $\mathrm{kcal} / \mathrm{mol}$ ). The other three abstraction reactions are thermoneutral and have barriers of ca. $12 \mathrm{kcal} / \mathrm{mol}$. Even though the structures for $\mathrm{BH}_{2} \mathrm{Cl}+\mathrm{H} \rightarrow \mathrm{BHCl}+\mathrm{H}_{2}$ and $\mathrm{BHCl}_{2}+\mathrm{H} \rightarrow$ $\mathrm{BCl}_{2}+\mathrm{H}_{2}$ are first-order saddle points, the structure for $\mathrm{BH}_{3}$


Figure 4. Hydrogen abstraction by H from $\mathrm{BH}, \mathrm{BH}_{3}, \mathrm{BHCl}$, and $\mathrm{BHCl}_{2}$ optimized at the MP2(full)/6-31G(d) level of theory. Note that $\mathrm{BH}_{2}-$ $\mathrm{H}-\mathrm{H}$ is a second-order saddle point, but all other structures are firstorder saddle points (transition states).


Figure 5. Intermediates for hydrogen addition to $\mathrm{BH}, \mathrm{BH}_{3}, \mathrm{BHCl}$, and $\mathrm{BHCl}_{2}$ optimized at the MP2(full)/6-31G(d) level of theory.
$+\mathrm{H} \rightarrow \mathrm{BH}_{2}+\mathrm{H}_{2}$ is a second-order saddle point (two imaginary frequencies).

Further optimization reveals that the $\mathrm{BH}_{3-n} \mathrm{Cl}_{n}+\mathrm{H} \rightarrow$ $\mathrm{BH}_{2-n} \mathrm{Cl}_{n}+\mathrm{H}_{2}$ reactions can proceed by a lower energy pathway involving H addition to form a $\mathrm{BH}_{4-n} C l_{n}$ intermediate followed by elimination of $\mathrm{H}_{2}$.

$\mathrm{BH}_{4-\mathrm{n}} \mathrm{Cl}_{\mathrm{n}}$ addition intermediate

The hydrogen atom approaches along the axis of the empty $p$ orbital on boron with little or no barrier to form a highly distorted tetracoordinate intermediate, as shown in Figure 5. The $\mathrm{BH}_{4}$ structure is isoelectronic with $\mathrm{CH}_{4}{ }^{+}$and $\mathrm{SiH}_{4}{ }^{+}$; each of these species suffers Jahn-Teller distortions ${ }^{27}$ because each can be obtained by removing an electron from a triply degenerate orbital in tetrahedral $\mathrm{BH}_{4}^{-}, \mathrm{CH}_{4}$, or $\mathrm{SiH}_{4}$. For $\mathrm{BH}_{4}$, the lowest energy structure contains a three-membered ring with elongated $\mathrm{B}-\mathrm{H}$ and $\mathrm{H}-\mathrm{H}$ bonds and is ca. $15 \mathrm{kcal} / \mathrm{mol}$ more stable than $\mathrm{BH}_{3}+\mathrm{H}$ or $\mathrm{BH}_{2}+\mathrm{H}_{2}$. The structures of $\mathrm{BHCl}-\mathrm{H}_{2}$ and $\mathrm{BCl}_{2}-$


Figure 6. Transition states for the insertion of $\mathrm{BH}_{2}, \mathrm{BHCl}$, and $\mathrm{BCl}_{2}$ into $\mathrm{H}_{2}$ and the elimination of $\mathrm{H}_{2}$ from the intermediates $\mathrm{BH}_{2}-\mathrm{H}_{2}$, $\mathrm{BHCl}-\mathrm{H}_{2}$, and $\mathrm{BCl}_{2}-\mathrm{H}_{2}$ optimized at the MP2(full)/6-31G(d) level of theory.


Figure 7. Transition states for hydrogen abstraction by Cl from BH , $\mathrm{BH}_{3}, \mathrm{BH}_{2} \mathrm{Cl}$, and $\mathrm{BHCl}_{2}$ optimized at the MP2(full)/6-31G(d) level of theory.
$\mathrm{H}_{2}$ are analogous but are only ca. 7 and $4 \mathrm{kcal} / \mathrm{mol}$ more stable than the separated species.

The elimination reaction $\mathrm{BH}_{4-n} \mathrm{Cl}_{n} \rightarrow \mathrm{BH}_{2-n} \mathrm{Cl}_{n}+\mathrm{H}_{2}$ proceed via the transition states shown in Figure 6. The reverse of these reactions are insertions and behave similar to carbene and silylene insertion reactions, ${ }^{26} \mathrm{AXY}+\mathrm{H}_{2} \rightarrow \mathrm{AH}_{2} \mathrm{XY}, \mathrm{A}=\mathrm{C}$, Si . The transition states are non-least motion and occur progressively later along the reaction path, and the insertion barriers increase with halogen substitution ( $0.4,6.7$, and 12.0 $\mathrm{kcal} / \mathrm{mol}$ for $\mathrm{BH}_{2}, \mathrm{BHCl}$, and $\mathrm{BCl}_{2}+\mathrm{H}_{2}$, respectively). For $\mathrm{BH}_{3}+\mathrm{H} \rightarrow \mathrm{BH}_{2}+\mathrm{H}_{2}$, the barrier for addition-elimination ( $1.1 \mathrm{kcal} / \mathrm{mol}$ ) is considerably lower than direct abstraction (12.5 $\mathrm{kcal} / \mathrm{mol}$, second-order saddle point). However, for $\mathrm{BHCl}_{2}+$ $\mathrm{H} \rightarrow \mathrm{BCl}_{2}+\mathrm{H}_{2}$, the two processes have nearly equal barriers and will be competitive ( $12.2 \mathrm{kcal} / \mathrm{mol}$ for addition-elimination vs $11.7 \mathrm{kcal} / \mathrm{mol}$ for abstraction).

For the divalent boron species, the $\mathrm{BH}_{2-n} \mathrm{Cl}_{n}+\mathrm{H} \rightarrow$ $\mathrm{BH}_{1-n} \mathrm{Cl}_{n}+\mathrm{H}_{2}$ process also occurs by addition-elimination. Even though the reverse of the elimination step can have a large barrier (Figure 2, $4.6 \mathrm{kcal} / \mathrm{mol}$ for $\mathrm{BH}+\mathrm{H}_{2}$, but $31.7 \mathrm{kcal} / \mathrm{mol}$ for $\mathrm{BCl}+\mathrm{H}_{2}$ ), the overall process is quit exothermic ( $27.1 \mathrm{kcal} /$ mol for $\mathrm{BH}_{2}+\mathrm{H} \rightarrow \mathrm{BH}+\mathrm{H}_{2}$ and $45.4 \mathrm{kcal} / \mathrm{mol}$ for $\mathrm{BHCl}+$ $\left.\mathrm{H} \rightarrow \mathrm{BCl}+\mathrm{H}_{2}\right)$ and there is no net barrier.

The transition states for the abstraction of H by Cl are shown in Figure 7. For $\mathrm{BH}+\mathrm{Cl} \rightarrow \mathrm{B}+\mathrm{HCl}$, the reaction is exothermic by $19.9 \mathrm{kcal} / \mathrm{mol}$ and appears to occur by a linear


Figure 8. Transition states for chlorine abstraction by H from BCl , $\mathrm{BH}_{2} \mathrm{Cl}, \mathrm{BHCl}_{2}$, and $\mathrm{BCl}_{3}$ optimized at the MP2(full)/6-31G(d) level of theory.
approach without a barrier. The $\mathrm{BH}_{2}$ and BHCl plus Cl reactions are exothermic ( 25.2 and $43.8 \mathrm{kcal} / \mathrm{mol}$, respectively) and proceed by barrierless addition to form $\mathrm{BH}_{2-n} \mathrm{Cl}_{n+1}$ which can eliminate HCl (see Figure 3), with no net barrier. The $\mathrm{BH}_{3-n} \mathrm{Cl}_{n}+\mathrm{Cl} \rightarrow \mathrm{BH}_{2-n} \mathrm{Cl}_{n}+\mathrm{HCl}$ reactions are nearly thermoneutral with little or no barrier. An addition-elimination mechanism is also conceivable but would involve more extensive rearrangement than the $\mathrm{BH}_{3-n} \mathrm{Cl}_{n}+\mathrm{H} \rightarrow \mathrm{BH}_{2-n} \mathrm{Cl}_{n}+\mathrm{H}_{2}$ case.

The transition states for $\mathrm{BCl}+\mathrm{H} \rightarrow \mathrm{B}+\mathrm{HCl}$ and $\mathrm{BH}_{3-n} \mathrm{Cl}_{n}$ $+\mathrm{H} \rightarrow \mathrm{BH}_{2-n} \mathrm{Cl}_{n}+\mathrm{HCl}$ are shown in Figure 8. In contrast to the abstraction reactions discussed above, these transition states have a bent $\mathrm{B}-\mathrm{Cl}-\mathrm{H}$ moiety. The reactions are endothermic by ca. $20 \mathrm{kcal} / \mathrm{mol}$; the reverse barriers are relatively small but increase with increasing chlorine substitution. Reaction path following confirms that these are abstraction reactions.

## Summary

Although the reactions in Figure 1 are numerous, they appear to involve only three simple classes: bond cleavage, radical abstraction, and insertion reactions. Closer examination, however, reveals a richer structure. Insertion reactions occur by highly distorted, non-least motion transition states. Insertion barriers increase with chlorine substitution on boron, whereas abstraction barriers are relatively insensitive. The transition states for the abstraction of hydrogen are linear, but those for the abstraction of chlorine are bent. Hydrogen abstraction by chlorine proceeds with little or no barrier, but hydrogen abstraction by hydrogen has a barrier of ca. $12 \mathrm{kcal} / \mathrm{mol}$. The lowest energy pathway for the latter, $\mathrm{BH}_{3-n} \mathrm{Cl}_{n}+\mathrm{H} \rightarrow$ $\mathrm{BH}_{2-n} \mathrm{Cl}_{n}+\mathrm{H}_{2}$, is not an abstraction but an additionelimination mechanism via a tetracoordinate intermediate. Likewise, the divalent species undergo addition-elimination reactions, rather than abstractions. All of the transition structures have been characterized at the MP2(full)/6-31G(d) level and the barriers have been computed at the G2 level of theory. These calculations should provide a foundation for a detailed analysis of reaction rates in the $\mathrm{BH}_{m} \mathrm{Cl}_{n}$ system.

Acknowledgment. This work was supported by grants from the Office of Naval Research (for S.J.H., N00014-95-C-0155)
and from the National Science Foundation (for H.B.S., CHE9400678).

## References and Notes

(1) Olsson, M.; Stridh, B.; Jansson, U.; Carlsson, J.; Soderberg, S. Mater. Sci. Eng. 1988, A105, 453. Gafri, O.; Grill, A.; Itzhak, D.; Inspektor, A.; Avni, R. Thin Solid Films 1980, 72, 523. Friesen, T.; Haupt, J.; Gissler, W.; Barna, A.; Barna, P. B. Surf. Coatings Technol. 1991, 48, 169.
(2) Rebenne, H.; Pollard, R. J. Am. Ceramic Soc. 1987, 70, 907.
(3) Schlegel, H. B.; Harris, S. J. J. Phys. Chem. 1994, 98, 11178.
(4) Rice, J. K.; Caldwell, N. J.; Nelson, H. H. J. Phys. Chem. 1989, 93, 3600.
(5) Caldwell, N. J.; Rice, J. K.; Nelson, H. H.; Adams, G. F.; Page, M. J. Chem. Phys. 1990, 93, 479.
(6) Jourdain, J. L.; Laverdet, G.; Le Bras, G.; Combourieu, J. J. Chim. Phys. 1981, 78, 253.
(7) Garrett, B. C.; Truhlar, D. G. J. Am. Chem. Soc. 1979, 101, 5207.
(8) Mayer, S. W.; Schieler, L.; Johnston, H. S. 11th Symposium (International) on Combustion, [Proceedings]; The Combustion Institute: Pittsburgh, PA, 1967; p 837.
(9) Curtiss, L. A.; Jones, C.; Trucks, G. W.; Raghavachari, K.; Pople, J. A. J. Chem. Phys. 1990, 93, 2537.
(10) Curtiss, L. A.; Raghavachari, K.; Trucks, G. W.; Pople, J. A. J. Chem. Phys. 1991, 94, 7221.
(11) Curtiss, L. A.; Carpenter, J. E.; Raghavachari, K.; Pople, J. A. J. Chem. Phys. 1992, 96, 9030.
(12) Frisch, M. J.; Trucks, G. W.; Schlegel, H. B.; Gill, P. M. W.; Johnson, B. G.; Robb, M. A.; Cheeseman, J. R.; Keith, T.; Petersson, G. A.; Montgomery, J. A.; Raghavachari, K.; Al-Laham, M. A.; Zakrzewski, V. G.; Ortiz, J. V.; Foresman, J. B.; Cioslowski, J.; Stefanov, B. B.; Nanayakkara, A.; Challacombe, M.; Peng, C. Y.; Ayala, P. Y.; Chen, W.; Wong, M. W.; Andres, J. L.; Replogle, E. S.; Gomperts, R.; Martin, R. L.; Fox, D. J.; Binkley, J. S.; Defrees, D. J.; Baker, J.; Stewart, J. P.; HeadGordon, M.; Gonzalez, C.; Pople, J. A. GAUSSIAN 94; Gaussian, Inc.: Pittsburgh PA, 1995.
(13) Hariharan, P. C.; Pople, J. A. Theor. Chim. Acta 1973, 28, 213 and references therein. Francl, M. M.; Pietro, W. J.; Hehre, W. J.; Binkley, J. S.; Gordon, M. S.; DeFrees, D. J.; Pople, J. A. J. Chem. Phys. 1982, 77, 3654. Frisch, M. J.; Pople, J. A.; Binkley, J. S. J. Chem. Phys. 1984, 80, 3265 and references therein. The $6-31 \mathrm{G}(\mathrm{d})$ and $6-31 \mathrm{G}(\mathrm{d}, \mathrm{p})$ basis sets are the same as $6-31 \mathrm{G}^{*}$ and $6-31 \mathrm{G}^{* *}$, respectively.
(14) Schlegel, H. B. J. Comput. Chem. 1982, 3, 214. Peng, C.; Ayala, P. Y.; Schlegel, H. B.; Frisch, M. J. J. Comput. Chem. 1996, 17, 49.
(15) Pople, J. A.; Krishnan, R.; Schlegel, H. B.; Binkley, J. S. Int. J. Quantum Chem., Quantum Chem. Symp. 1979, 13, 225.
(16) Trucks, G. W.; Frisch, M. J.; Andres, J. L.; Schlegel, H. B. J. Chem. Phys., submitted for publication.
(17) Gonzalez, C.; Schlegel, H. B. J. Chem. Phys. 1989, 90, 2154. Gonzalez, C.; Schlegel, H. B. J. Phys. Chem. 1990, 94, 5523.
(18) Møller, C.; Plesset, M. S. Phys. Rev. 1934, 46, 618. For a review: Bartlett, R. J. Annu. Rev. Phys. Chem. 1981, 32, 359.
(19) Pople, J. A.; Head-Gordon, M.; Raghavachari, K. J. Chem. Phys. 1987, 87, 5968.
(20) Pople, J. A.; Scott, A. P.; Wong, M. W.; Radom, L. Isr. J. Chem. 1993, 33, 345 .
(21) Pople, J. A.; Schlegel, H. B.; Krishnan, R.; DeFrees, D. J.; Binkley, J. S.; Frisch, M. J.; Whiteside, R. A.; Hout, R. F.; Hehre, W. J. Int. J. Quantum. Chem., Quantum Chem. Symp. 1981, 15, 269.
(22) Storms, E.; Mueller, B. J. Phys. Chem. 1977, 81, 318.
(23) Ruscic, B.; Mayhew, C. A.; Berkowitz, J. J. Chem. Phys. 1988, 88, 5580.
(24) Ochterski, J. W.; Petersson, G. A.; Wiberg, K. B. J. Am. Chem. Soc. 1995, 117, 11299.
(25) Chase, M. W.; Davies, C. A.; Downey, J. R.; Frurip, D. J.; McDonald, R. A.; Szverud, A. N. JANAF Thermochemical Tables 3rd ed.; J. Phys. Chem. Ref. Data 1985, 14.
(26) Sosa, C.; Schlegel, H. B. J. Am. Chem. Soc. 1984, 106, 5847. Ignacio, E. W.; Schlegel, H. B. J. Phys. Chem. 1992, 96, 1620. Su, M.-D.; Schlegel, H. B. J. Phys. Chem. 1993, 97, 9981.
(27) For leading references see: Frey, R. F.; Davidson, E. R. J. Chem. Phys. 1988, 88, 1775; 1988, 89, 4227.

JP960347H


[^0]:    ${ }^{\otimes}$ Abstract published in Advance ACS Abstracts, May 1, 1996.

